

# On the relationship between glottal pulse shape and its spectrum: correlations of open quotient, pulse skew and peak flow with source harmonic amplitudes

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## Abstract

This paper explores the relationship between the glottal pulse amplitude ( $U_p$ ) and the amplitude of the first harmonic ( $H1$ ), as well as the combined effects of  $U_p$ , the open quotient ( $O_q$ ) and degree of pulse asymmetry/skew ( $R_k$ ) on the low end of the source spectrum. This serves to elucidate their relationship to the  $H1$ - $H2$  estimate, widely used to make inferences on changes in  $O_q$  and voice quality. It has been suggested that  $H1$  is mainly determined by  $U_p$  and that the pulse shape has a relatively small impact. To investigate this, a series of glottal pulses were generated using the LF model, where  $U_p$  was kept constant, while  $O_q$  and  $R_k$  were systematically varied. The resulting harmonic amplitudes of these pulses show that  $U_p$  is not the sole determinant of  $H1$ . Rather,  $H1$  is highly dependent on  $O_q$  and to a certain degree also on  $R_k$ . Although the effects of these parameters on the lowest harmonics is rather complex, we find that the  $H1$ - $H2$  measure is broadly correlated with  $O_q$ . However, there is also a strong effect of differences in glottal skew, particularly at high  $O_q$  values, which could invalidate inferences on  $O_q$  and voice quality from estimates of  $H1$ - $H2$ .

**Index Terms:** voice source, glottal flow, LF model, open quotient, skew,  $H1$ - $H2$ , frequency domain

## 1. Introduction

The present study is part of an ongoing investigation of the detailed mapping of time and frequency domain dimensions of the voice source, important to our understanding of voice source variation in speech. Each domain yields different insights: whereas the time domain illuminates production, the frequency dimension relates more directly to perception.

Our past research has entailed detailed analyses of the voice source, using inverse filtering techniques and the modelling of the glottal waveform [1] using the Liljencrants-Fant (LF) model [2]. These have examined the source correlates of voice qualities [3, 4], the variation of the voice in prosody, to signal linguistic [5-7] and affective information [8-11].

Such time domain voice source analyses are necessarily limited in scope due in part to the fact that they require stringent recording conditions, e.g., to ensure phase-linearity (for a fuller discussion of the limiting factors, see [1]). Consequently, there is a relatively limited body of analytic data available, and researchers have tended to use spectral information as proxy measures of the source. Such a measure is  $H1$ - $H2$ , which is widely taken to be an indicator of the open quotient ( $O_q$ ) of the glottal pulse [12], and long used as a measure of the breathy-tense dimension of voice quality [13-18]. An increase in  $O_q$  is assumed to lead to an increase in the

amplitude of  $H1$ , thus increasing  $H1$ - $H2$  [19]. (Note that when the measure is based on the speech waveform, the  $H1^*$ - $H2^*$  measure is often used, to correct for the vocal tract resonances [20, 12].)

Although a full time-frequency mapping is our objective, the focus of this paper is on the lower end of the source spectrum. We explore the relationship of the amplitude of the glottal flow pulse (i.e. the peak flow,  $U_p$ ) and the amplitude of the first harmonic of the source spectrum ( $H1$ ). Furthermore, we examine the combined effects of  $U_p$ , the open quotient ( $O_q$ ) and pulse asymmetry/skew ( $R_k$ ) on the low end of the source spectrum, looking in detail at how  $H1$ ,  $H2$ ,  $H3$  and the  $H1$ - $H2$  estimate relate to  $O_q$ .

Fant and Lin [21] provide a theoretical basis for the relationship between time and frequency domain representations of the voice source and they show that the amplitude of the first harmonic of the source spectrum can be derived according to (1). Here the effect of the vocal tract filter is disregarded (or assumed to have been properly eliminated by inverse filtering) and  $|R(f)|$  represents the effect of the radiation transfer on the amplitude spectrum of the speech output at a specific distance from the lip opening of the speaker.

$$H1 = k \cdot \frac{U_p}{2} |R(f)| \quad (1)$$

If we ignore the radiation factor, we end up with  $H1$  being determined by  $k \cdot U_p/2$ , i.e. half the peak flow scaled by the value of  $k$ , which is a correction factor. In this case  $H1$  would correspond to the amplitude of the first harmonic of the waveform produced at the glottis, and that is how the term  $H1$  is used in the following text.

As pointed out by Fant and Lin [21], for approximate calculations, it is convenient to set the correction factor  $k = 1$ . Although it is explained that  $k$  does vary with  $O_q$ , no detailed information is presented on precisely how the glottal pulse shape affects  $k$  (and  $H1$ ).

It is hoped that a growing understanding of the time-frequency mapping of source dimensions will facilitate the development of more robust analysis techniques, not constrained by factors such as the recording conditions. We further see it as key to establishing how source dimensions map to the auditory perception of voice quality (see for example [22]), as these are not well understood. Ultimately the goal is to understand and model how we use modulation of the voice for the rich nuancing of prosody in speech communication.

In a parallel research strand, we are developing text-to-speech synthesis for the dialects of Irish (see [23, 24]), and plan to exploit them in educational games [25] and dialogue systems [26]. Our basic research on voice modelling thus goes

hand in hand with our goal to provide more adequate descriptions of voice modulation in prosody, which is essential for this kind of application where speaker affect is important.

## 2. Methods

To analyse the relationships between the glottal pulse shape and the amplitude of source spectral components, the LF model [2] was used to generate glottal pulses with constant peak amplitude ( $U_p$ ) and pulse duration, while the open quotient and the degree of pulse asymmetry were varied in controlled steps by changing the  $O_q$  and  $R_k$  parameters, as defined in Figure 1.

The LF model (see Figure 1) is determined by the two expressions in (2), which generate the differentiated glottal waveform of the open phase and return phase respectively. Note that as our focus here is on  $O_q$  and  $R_k$ , the pulses generated did not have a return phase (i.e.  $T_a = 0$ ).

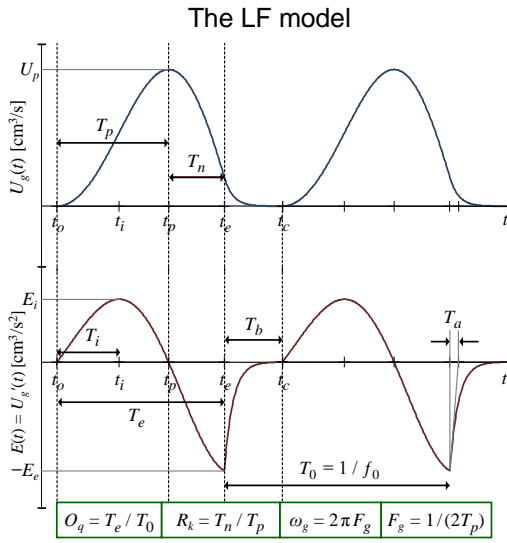


Figure 1. Two LF model pulses and parameter definitions (for details, see [1]). Glottal flow (top), flow derivative (bottom).

$$U_g'(t) = \begin{cases} E_0 e^{\alpha t} \sin \omega_g t & t_o \leq t \leq t_e \\ \frac{-E_e}{\varepsilon T_a} \left( e^{-\varepsilon(t-t_e)} - e^{-\varepsilon T_b} \right) & t_e < t < t_c \end{cases} \quad (2)$$

In this work we are particularly interested in the flow pulse, rather than the flow derivative, and therefore the alternative expressions shown in (3) are used, which are obtained by the integration of (2). They produce the open phase and return phase of the LF glottal flow pulse (see also [27, 28]).

As pointed out in [29], the actual LF model parameters are  $E_e$ ,  $T_e$ ,  $\omega_g$ ,  $\alpha$ ,  $\varepsilon$  and  $T_b$ .  $E_0$  and  $T_a$  in (2) and (3) are only auxiliary parameters:  $E_0$  is determined by  $E_e$ ,  $\alpha$ ,  $T_e$ , and  $\omega_g$ , while  $T_a$  is determined by  $\varepsilon$  and  $T_b$ .

$E_e$  is the excitation strength,  $T_e$  is the duration of the open phase which determines  $O_q$ ,  $\omega_g$  is the angular frequency ( $2\pi F_g$ ) of the glottal pulse in the open phase and  $\alpha$  relates to its bandwidth. Apart from these parameters defining the LF

pulse, the model requires that there is area balance [2], i.e. that the area of the positive part of the differentiated glottal pulse equals the area of the negative part (see the lower panel of Figure 1), so as to avoid any drift in the zero-flow line. To achieve this,  $\alpha$  is allowed to ‘float’, i.e. it is implicitly determined by the other parameters in order to obtain area balance.

It is clear from the above that  $U_p$  is not a parameter of the LF model. Therefore, as for  $\alpha$ ,  $U_p$  is typically allowed to float, and will attain a value entirely determined by the settings of the other parameters.

$$U_g(t) = \begin{cases} \frac{E_0 e^{\alpha t} \sin \left( \omega_g t - \arctan \frac{\omega_g}{\alpha} \right)}{\sqrt{\alpha^2 + \omega_g^2}} + \frac{E_0 \omega_g}{\alpha^2 + \omega_g^2} & t_o \leq t \leq t_e \\ \frac{E_e}{\varepsilon^2 T_a} \left[ e^{-\varepsilon(t-t_e)} + \varepsilon e^{-\varepsilon T_b} \left( t - \left( t_c + \frac{1}{\varepsilon} \right) \right) \right] & t_e < t < t_c \end{cases} \quad (3)$$

This does not serve the present purpose where we wish to keep  $U_p$  constant while varying  $O_q$  and  $R_k$ . To do this, we use a version of the iterative procedure described in [29], where the direct control of  $\omega_g$  is sacrificed, so that  $U_p$  can be directly controlled.

However, we also wish to control the glottal skew by varying  $R_k$  in a systematic way, and since  $\omega_g$  is allowed to float, this also means that  $R_k$  cannot be controlled (see Figure 1). Therefore, the algorithm in [29] was modified so that  $E_e$  would float instead of  $\omega_g$ . This was achieved by replacing the two partial derivatives with respect to  $\omega_g$  in [29] with the following partial derivatives with respect to  $E_e$ , shown in (4) and (5).

$$\frac{\partial f_1}{\partial E_e} = \pi \omega_g \left( e^{-\alpha T_e} + e^{-\alpha (T_e - \pi \omega_g^{-1})} \right) \quad (4)$$

$$\frac{\partial f_2}{\partial E_e} = 0 \quad (5)$$

Using this modified parameter control of the LF model, 81 different flow pulses were generated.  $U_p$  remained fixed and was arbitrarily set to 10. The glottal period (i.e.  $T_e + T_b$ , see Figure 1) was kept constant at 10 ms.

Nine different  $O_q$  settings were used, ranging from 0.15 to 0.95 in steps of 0.1, covering a very wide range of  $O_q$  values. For each of the  $O_q$  settings, nine pulses with different  $R_k$  values were produced, also ranging from 0.15 to 0.95 in steps of 0.1 – again, covering most of the possible range of  $R_k$  values.

The sampling frequency was 20 kHz, which was deemed sufficiently high to avoid any impact of aliasing on the lower end of the source spectrum. Each pulse was repeated five times in order to produce a harmonic spectrum. A 1000-point (50 ms, rectangular window) DFT spectrum was calculated for each of the 81 glottal waveforms, and the amplitudes of the first three harmonics were extracted. The window size was chosen so as to ensure that the output frequency samples would coincide with the harmonic frequencies, thus avoiding potential rounding errors.

### 3. Results of Spectral Analysis

#### 3.1. Influence of $O_q$ and $R_k$ on $H1$ , $H2$ and $H3$

Figure 2 shows how the amplitudes (in dB) of the first three harmonics vary as a function of  $O_q$ . Each panel shows this variation for different  $R_k$  settings, 0.15, 0.35 and 0.65 in the top, mid and lower panels respectively. Figure 3 shows in separate panels the variation in  $H1$  and  $H2$  (on a linear scale) with  $O_q$ , with the nine different  $R_k$  settings superimposed. 0 dB corresponds to  $k = 1$ , which here is the (linear) value of 5.

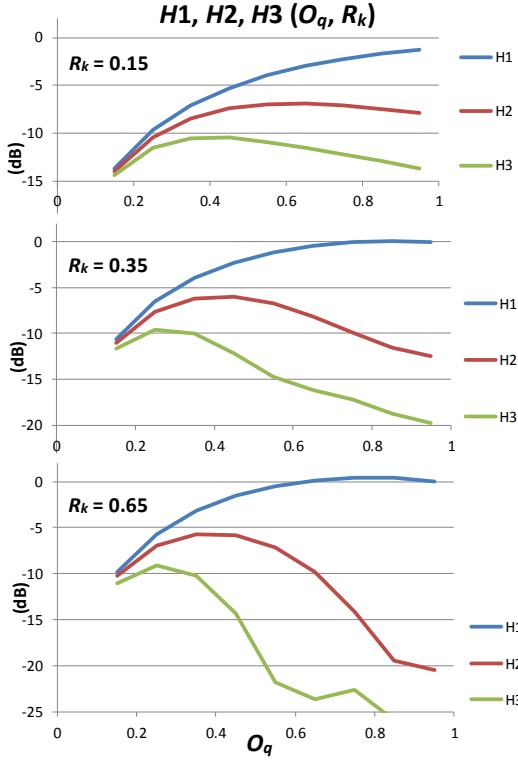


Figure 2:  $H1$ ,  $H2$  and  $H3$  amplitude levels as a function of  $O_q$  for three different  $R_k$  settings.

Figures 2 and 3 show that  $H1$  varies considerably as a function of  $O_q$  and that there is also a substantial influence of  $R_k$ . At relatively low  $O_q$  values  $H1$  increases with increasing  $O_q$ . However, this increase plateaus and  $H1$  even drops (Figure 3) at very high  $O_q$  (except where  $R_k$  is very low). Once the plateau is reached, the further changes in  $H1$  are relatively minor.

It is also clear in the upper panel of Figure 3 that  $R_k$  has a strong influence on  $H1$ , particularly when the pulse is skewed, in the range of 0.15 to 0.35. The influence of  $R_k$  extends up to about 0.45.

The lower panel of Figure 3 shows how  $H2$  is influenced by the pulse shape parameters. At very low  $O_q$  values  $H2$  rises with  $O_q$ , but drops thereafter, as  $O_q$  continues to rise. The  $O_q$  value where this drop occurs and its steepness depends on the  $R_k$  setting, being more extreme with high  $R_k$  (low skew).  $H3$  also drops sharply with rising  $O_q$  values, and also shows a strong influence of  $R_k$ .

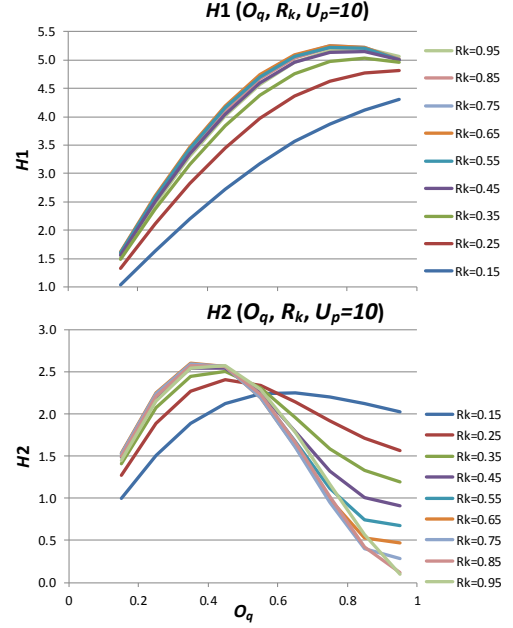


Figure 3:  $H1$  and  $H2$  as a function of  $O_q$  and  $R_k$ .  $U_p = 10$  in all cases.

#### 3.2. $H1$ - $H2$ as a measure of $O_q$

As mentioned in the Introduction, the  $H1$ - $H2$  measure has been widely interpreted in the literature as an indirect measure of  $O_q$ . However, it is clear from Figures 2 and 3 that neither  $H1$  nor  $H2$  have a straightforward correlation with  $O_q$ , even if  $R_k$  is held constant. Nonetheless, the difference between them is more closely correlated to  $O_q$ . And although the  $H1$ - $H2$  difference is generally assumed to be the consequence of a rise in  $H1$  with  $O_q$ , Figure 2 suggests that the drop in  $H2$  is likely to be an important factor.

A comparison of the three panels of Figure 2, shows that, despite a broad, positive correlation of  $O_q$  with  $H1$ - $H2$ , this correlation is greatly affected by the pulse skew  $R_k$ . This last is elaborated in detail in Figure 4, which charts how the  $H1$ - $H2$  correlation with  $O_q$  is impacted by variation in pulse skew. As typical values of  $R_k$  tend to range from about 0.2 to 0.6, for much of running speech, the  $H1$ - $H2$  measure is neither a direct nor unique indicator of the open quotient. This supports suggestions by [30] that  $R_k$  is likely to have an important impact on the lower part of the source spectrum and on the  $H1$ - $H2$  measure.

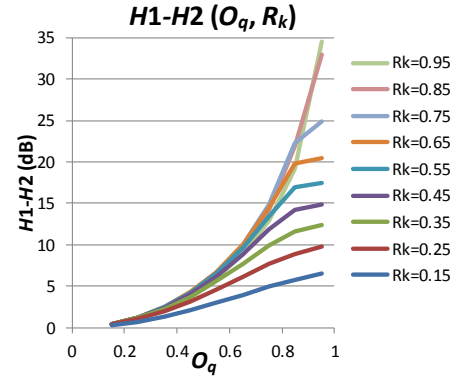


Figure 4: The  $H1$ - $H2$  as a function of  $O_q$  and  $R_k$ .

To sum up, while there are clear correspondences between the spectral measure  $H1$ - $H2$  and the open quotient, given the clear impact of other factors (particularly here, the pulse skew), one should be cautious in making inferences on  $O_q$  on the basis of  $H1$ - $H2$ . Though it may capture broad trends much of the time, its potential limitations need to be borne in mind.

The present analysis has been limited to the joint influences of  $U_p$ ,  $O_q$  and  $R_k$ , as these are considered to be the primary determinants of the lower end of the source spectrum. The analysis does not extend to consideration of other possible influences, such as variation of the return phase ( $T_a$ , see Figure 1), which primarily determines the shaping of the upper end of the source spectrum. Nonetheless, interactions are likely, and these will need to be considered at a later point.

#### 4. Predicting $H1$ and deriving the $k$ correction factor

We use these spectral measurements to predict more precisely the  $k$  factor. The variation in  $H1$  due to changes in  $O_q$ , closely follows part of a parabolic curve (see Figure 3). Second order polynomial fitting yields  $R^2$  values close to 1 for the nine  $R_k$  values ( $R^2 = 0.9992$  or higher). Thus, we can use a quadratic function to derive an estimate of the correction factor  $k$  according to equation (6), and, given the amplitude of the peak flow ( $U_p$ ),  $H1$  can be approximated according to equation (1) in the Introduction.

$$k = a_2 O_q^2 + a_1 O_q + a_0 \quad (6)$$

However, the coefficients of the function in (6) depend on  $R_k$ . The variation in the three  $a$ -coefficients of the polynomial in (6) also closely match quadratic functions (see Figure 5, left panel). Although the fit is very good, the  $R^2$  values were somewhat lower in this case (but not lower than 0.975). The three equation for the  $a$ -coefficients as a function of  $R_k$  are shown in (7), and the numerical values of the nine coefficients are shown in Table 1. The values have been rounded to three-digit precision, with negligible effect on the accuracy of predictions.

$$\begin{aligned} a_2 &= b_2 R_k^2 + b_1 R_k + b_0 \\ a_1 &= c_2 R_k^2 + c_1 R_k + c_0 \\ a_0 &= d_2 R_k^2 + d_1 R_k + d_0 \end{aligned} \quad (7)$$

Table 1: The values of the coefficients in equation (7) used to derive the coefficients for the quadratic equation (6) estimating the correction factor,  $k$ .

|       |        |        |
|-------|--------|--------|
| $b_2$ | $b_1$  | $b_0$  |
| 3.95  | -5.42  | 0.0307 |
| $c_2$ | $c_1$  | $c_0$  |
| -4.35 | 6.01   | 0.811  |
| $d_2$ | $d_1$  | $d_0$  |
| 0.161 | -0.243 | 0.0203 |

To test the accuracy of the of  $k$  correction factor as determined by (6) and (7), the  $H1$  prediction errors were calculated on a different set of 64 LF pulses, where both  $O_q$  and  $R_k$  varied from 0.2 to 0.9 in steps of 0.1. As can be ascertained in the right panel of Figure 5, there is a high correlation between the estimated and actual  $H1$  amplitudes. The average error in the  $H1$  estimate was 0.14 dB with a maximum error of 0.35 dB.

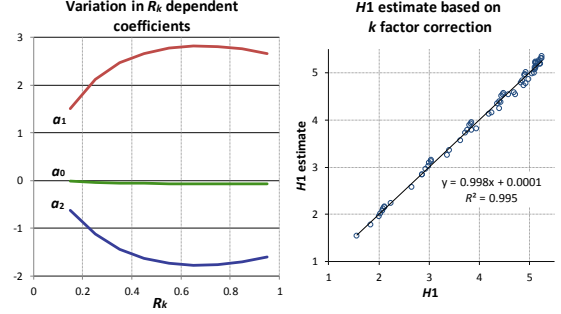


Figure 4: Variation in the  $R_k$  dependent  $a_2$ ,  $a_1$  and  $a_0$  coefficients of (6) (left panel).  $H1$  estimates vs. actual  $H1$  amplitudes using the  $k$  correction factor as defined by equation (6) (right panel).

#### 5. Conclusions

The approach taken in this paper – using adjustments to the LF model, to allow control of different source parameters – has permitted a detailed exploration of the influence of  $U_p$ ,  $O_q$  and  $R_k$  on the shaping of the low end of the source spectrum, and has illustrated the complex interaction of these parameters. A formula is presented that allows a precise estimation of the  $k$  factor and prediction of  $H1$ .

The spectral analyses also illuminate how these parameters affect  $H2$ , and the impact on the  $H1$ - $H2$  measure. Despite the broad correspondence with  $O_q$ , glottal skew was found to have a strong effect, particularly at high  $O_q$  values. These effects could potentially invalidate direct inferences from  $H1$ - $H2$  on  $O_q$  and on voice quality.

This work is seen as part of a larger enterprise of mapping time and frequency dimensions of the source. An extension of the present study will explore how the return phase of the glottal pulse, usually seen as mainly influencing the higher end of the source spectrum, impacts on the low frequency end.

The quadratic correspondence found between  $U_p$  and  $H1$  as the glottal pulse shape varies, means that there is not a one-to-one mapping: the same harmonic amplitudes can be produced by different glottal pulse configurations. Although the basic framework of Fant and Lin [21] is very useful, the more detailed elaboration of the precise correlations here indicate that the inverse prediction of time domain measures from the frequency domain will be complex and challenging. Future work will explore the potential of deriving the peak glottal flow from the low harmonics, by exploiting some of the correlations elaborated here to impose constraints on possible solutions.

Progress towards comprehensive time-frequency mapping would greatly further our ability to explore the perception of voice source variation. It would facilitate more flexible and robust voice source analysis [31], allowing us greater access to aspects of speech communication that are little understood. This we believe to be vital, if speech synthesis [32, 33] and related technologies are to be adequate for many applications.

#### 6. Acknowledgements

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